

Isoparallel problems as generalized nonholonomic systems

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XXII International Workshop on Differential Geometric Methods
in Theoretical Physics

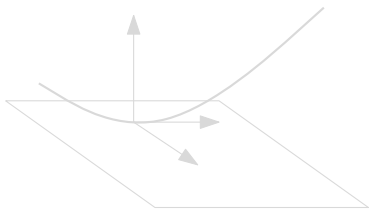
Będlewo, August 25th, 2007

Lagrange–d'Alembert's principle

Constrained mechanical system

- $L: TQ \rightarrow \mathbb{R}$
- $C \subseteq TQ$ (distribution)

Principle of virtual work: No work is done by the constraint forces in any virtual displacement consistent with the constraints



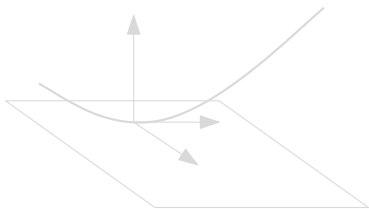
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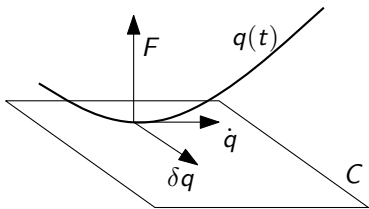
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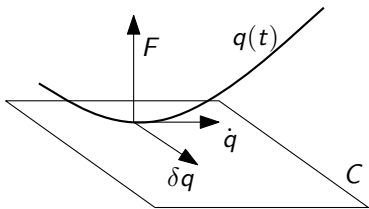
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$$F \in C^\circ \subset T^*Q$$

Lagrange–d'Alembert's principle

In a Lagrangian system,

$$F = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q}$$

so $F \in C^0$ means

$$\left(\frac{d}{dt} \frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} \right) \cdot \delta q = 0 \quad \text{for } \delta q \in C$$

(in addition, $\dot{q} \in C$)

$$F \in C^0 \quad \Rightarrow \quad F \cdot \dot{q} = 0 \text{ (power)} \quad \Rightarrow \quad \text{preservation of energy}$$

Constraints on motion + pple. of virtual work



constraints on variations

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- It is convenient
- Applies to a wide range of problems
- Covariance
- Reduction, discretization...
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Suitable generalization of Lagrange–d'Alembert's principle

- elastic rolling bodies, pneumatic tires
- dissipative systems
- servomechanisms, control strategies
- isoparallel problems

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Generalized nonholonomic systems

Generalized nonholonomic system (GNHS) on Q :

- $L: TQ \rightarrow \mathbb{R}$
- $C_K \subseteq TQ$ (kinematic distribution)
- $C_V \subseteq TQ$ (variational distribution)

A trajectory of a GNHS is a curve $q(t)$ such that

- $\dot{q} \in C_K$
- $\delta \int L(q, \dot{q}) dt = 0$ with respect to $\delta q \in C_V$, which implies

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Constraint forces: $F \in C_V^\circ$

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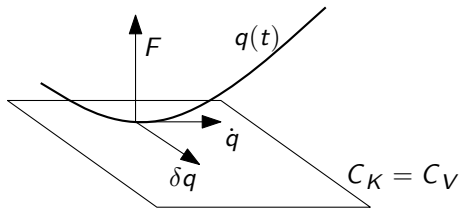
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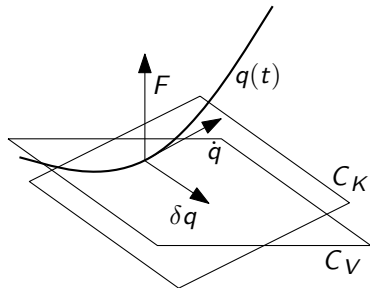
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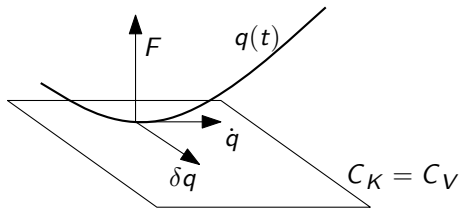
Nonholonomic
(Lagrange–d'Alembert)



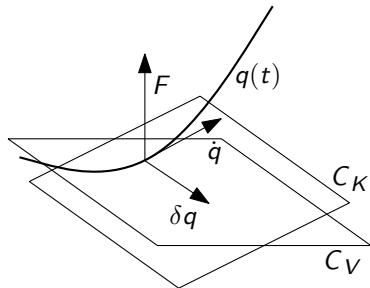
Generalized nonholonomic

Energy is not preserved in GNHS

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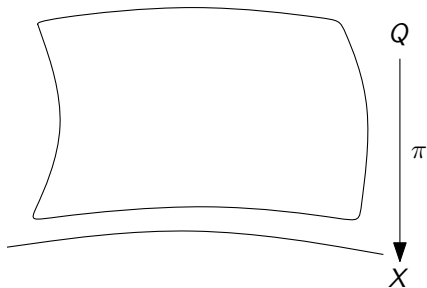


Generalized nonholonomic

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Isoparallel problems

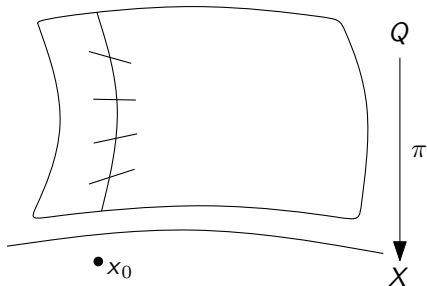
- Principal bundle Q , with structure group G
- Principal connection $A: TQ \rightarrow \mathfrak{g}$



- For each curve $x(t)$ joining x_0 to x_1 , there is a *parallel transport operator* mapping $\pi^{-1}(x_0)$ into $\pi^{-1}(x_1)$
- Many curves might share the same parallel transport operator

Isoparallel problems

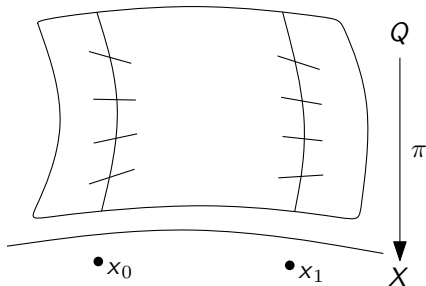
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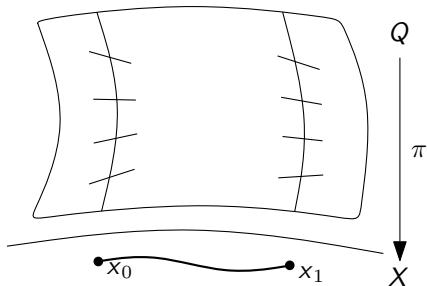
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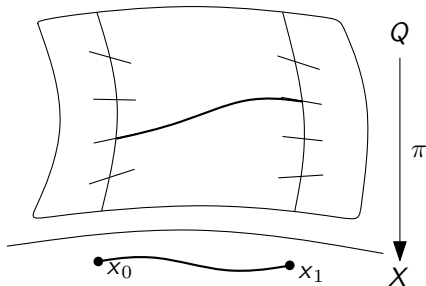
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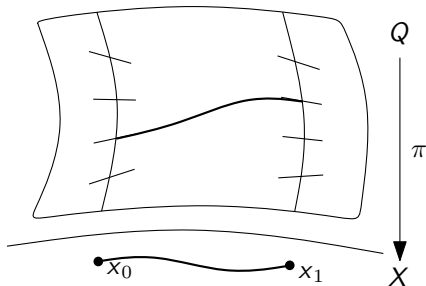
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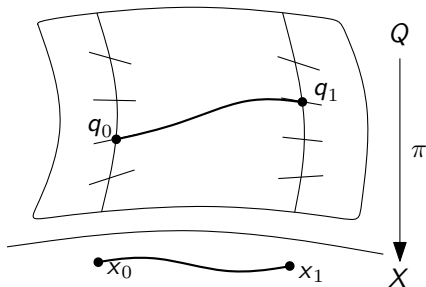
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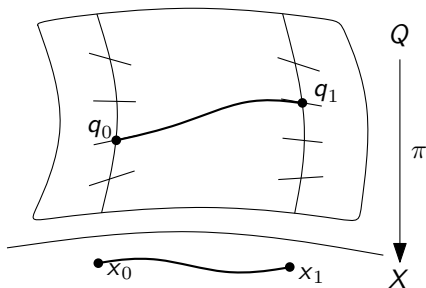
- Riemannian metric on the base

Isoparallel problem: Find a curve on X that extremizes length among those with a given parallel transport operator.

Equivalently: Among those horizontal curves joining q_0 to q_1 , find one whose projection extremizes length.

When $x_0 = x_1$, this is called the *isoholonomic* problem.

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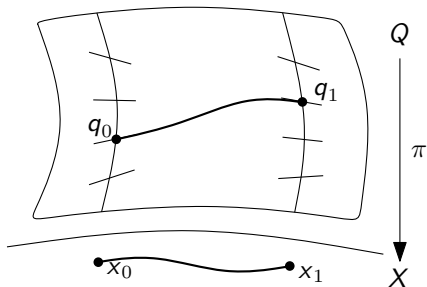
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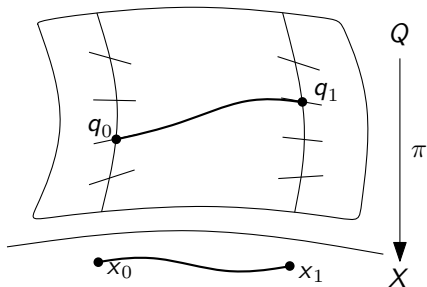
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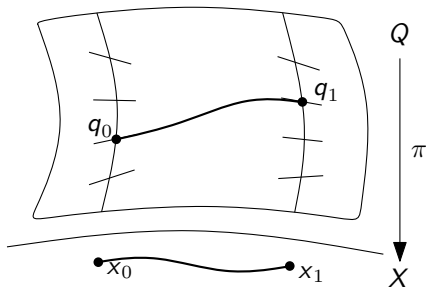
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For a given parallel transport operator...

- Is there any curve realizing it?
- Are there “enough” curves?
 - Working with C^1 curves, not always (*rigid curves*)
 - Working with absolutely continuous curves, yes (if horizontal distribution is bracket-generating)

Absolutely continuous curves provide a nice setting:

- Horizontal lifts are still well defined
- Mixed partial derivatives are equal in L^2
- Integration by parts works

We will assume the regular case

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Falling cat:

Q = space of positions of the cat as a deformable body

$G = \text{SO}(3)$

$X = Q/G$ = shapes (disregarding rigid rotations)

Horizontal distribution in TQ defined by “angular momentum equals zero”

Loop on X \dashrightarrow Hor. lift \dashrightarrow Reorientation

Metric on X measuring energy expenditure

► What is the most efficient way that the cat could change its shape so as to land on its feet?

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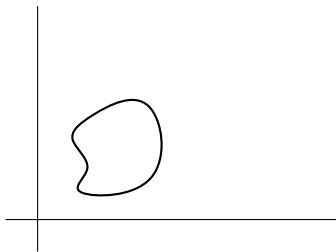
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Locomotion of a microorganism in a fluid



$Q \subset$ mappings from S^1 into \mathbb{R}^2

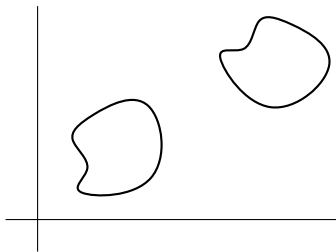
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$v \in TQ$: vector field on the membrane \dashrightarrow response of the fluid

Define a principal connection on Q and a metric on X

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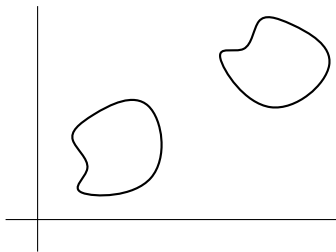
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Montgomery's theorem

Take a bi-invariant metric β on G (if there is any!)

If k is the metric on the base, define $L: TQ \rightarrow \mathbb{R}$ by

$$L(q, \dot{q}) = \frac{1}{2} \left(k(T\pi(q, \dot{q}), T\pi(q, \dot{q})) + \beta(A(q, \dot{q}), A(q, \dot{q})) \right)$$

(related to the metric $k \oplus \beta$ on Q)

► Then the (normal) solutions of the isoparallel problem are precisely the projections of the geodesics [Montgomery, 1990]

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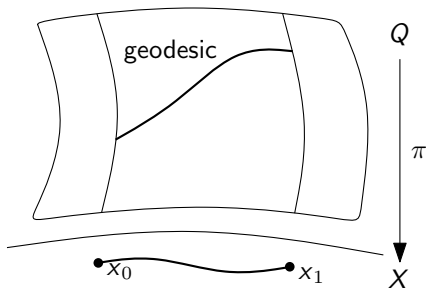
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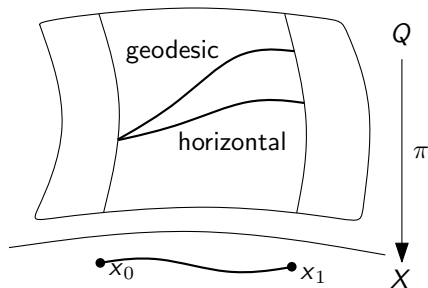


Lagrangian system \dashrightarrow project (and lift)

But... some groups do not admit a bi-invariant metric (e.g. $SE(2)$)

► Replace Lagrangian system by *generalized nonholonomic system*
[Cendra and Ferraro, 2006]

Montgomery's theorem

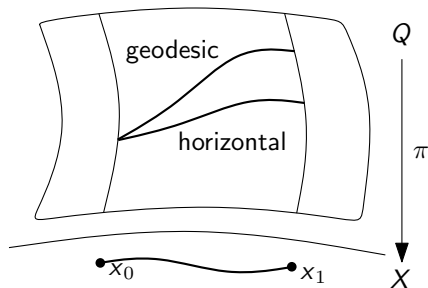


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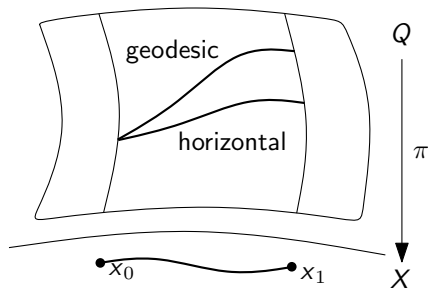


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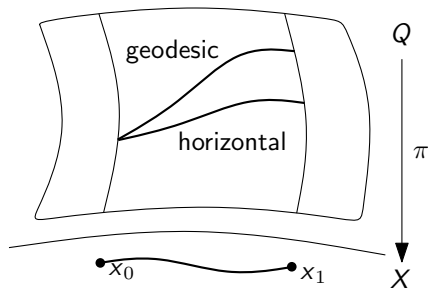


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$M =$ symmetric bilinear forms on \mathfrak{g} .

G -action: For $g \in G$, $\beta \in M$ and $\eta_1, \eta_2 \in \mathfrak{g}$, define

$$(g\beta)(\eta_1, \eta_2) = \beta(\text{Ad}_{g^{-1}} \eta_1, \text{Ad}_{g^{-1}} \eta_2).$$

Infinitesimal generator: If $\xi \in \mathfrak{g}$ and $\beta \in M$ then

$$(\xi\beta)(\eta_1, \eta_2) = \beta(-[\xi, \eta_1], \eta_2) + \beta(\eta_1, -[\xi, \eta_2]).$$

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Generalized nonholonomic system:

- It is a system on $Q \times M$
- $L: T(Q \times M) \rightarrow \mathbb{R}$,

$$L(q, \beta, \dot{q}, \dot{\beta}) = \frac{1}{2}k (T\pi(q, \dot{q}), T\pi(q, \dot{q})) + \frac{1}{2}\beta (A(q, \dot{q}), A(q, \dot{q}))$$

- Distributions in $T(Q \times M)$
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Associated bundles:

$$\tilde{M} = (Q \times M)/G$$

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These are bundles over $X = Q/G$ with standard fiber M (resp. \mathfrak{g}).

Notation:

$$\bar{\beta} = [q, \beta]_G \in \tilde{M}$$

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Reduced tangent bundle

$$T(Q \times M) \equiv TQ \times TM \equiv TQ \oplus (Q \times M) \oplus (Q \times M)$$

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The reduced Lagrangian $\ell: TX \oplus \tilde{\mathfrak{g}} \oplus 2\tilde{M} \rightarrow \mathbb{R}$ is

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Reduced variations

$$\begin{aligned}\delta x &= T\pi(\delta q) \\ \delta^A \bar{v} &= \tilde{B}(x)(\delta x, \dot{x}) + \frac{D\bar{\eta}}{Dt} + [\bar{v}, \bar{\eta}] \\ \delta^A \bar{\beta} &= -\bar{\eta}\bar{\beta}\end{aligned}$$

Reduced equations for the GNHS

$$\begin{aligned}\frac{D\bar{\beta}}{Dt} &= 0 \\ \frac{D\bar{v}}{Dt} &= 0 \\ (\nabla_{\dot{x}} \dot{x})^b &= -\bar{\beta}(\bar{v}, \tilde{B}(x)(\dot{x}, \cdot))\end{aligned}$$

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- Extremizing $\int_{t_0}^{t_1} \sqrt{k(T\pi(\dot{q}), T\pi(\dot{q}))} dt$
- Constrained to $A(q, \dot{q}) = 0$
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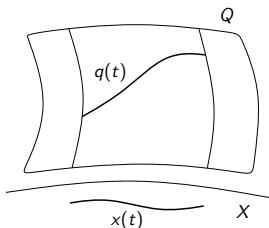
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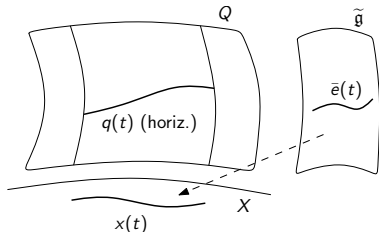


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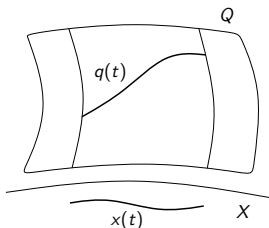


Reduced versions match, unreduced versions do not

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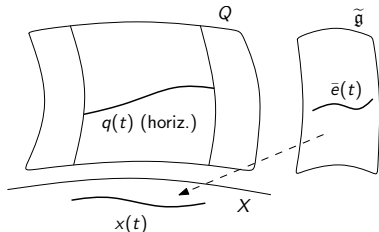


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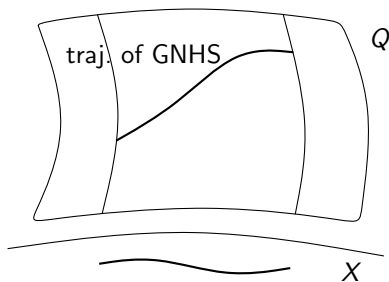
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Relationship with theorem on geodesics

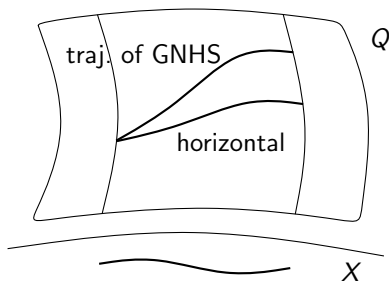


Trajectory of GNHS \dashrightarrow project and lift

Geodesic \dashrightarrow project and lift

With a bi-invariant metric, $C_K = TQ \oplus 0$ and $C_V = TQ \oplus 0$, and the GNHS becomes Montgomery's Lagrangian system

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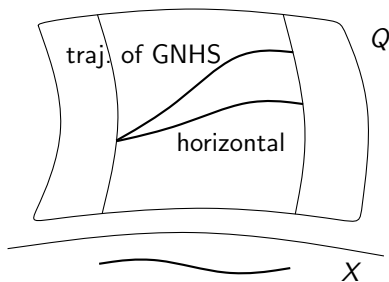


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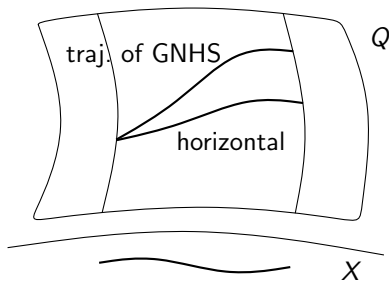


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 - Dissipative systems, control strategies
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- Disadvantage: no universal procedure to find C_V
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