

Symmetries and Quantum Entanglement

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Motivations

- Quantum Information Theory
- Quantum Computing
- Quantum Teleportation
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- structure of composite states in QM !!!

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Quantum entanglement — pure states

$$\mathcal{H}_{\text{total}} = \mathcal{H}_A \otimes \mathcal{H}_B$$

Definition:

$\psi \in \mathcal{H}_{\text{total}}$ is separable iff

$$\psi = \psi_A \otimes \psi_B$$

with $\psi_A \in \mathcal{H}_A$ & $\psi_B \in \mathcal{H}_B$.

Definition:

$\psi \in \mathcal{H}_{\text{total}}$ is entangled (inseparable) iff it is NOT separable.

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Test

$$\psi \otimes \varphi ?$$

$$\psi \otimes \chi + \varphi \otimes \omega + \psi \otimes \omega + \varphi \otimes \chi ?$$

$$= (\psi + \varphi) \otimes (\chi + \omega) !!!$$

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How to check for separability ???

Quantum entanglement — pure states

Schmidt decomposition

$$\psi \in \mathcal{H}_{\text{total}}$$

$$e_1, e_2, \dots \in \mathcal{H}_A \quad \& \quad f_1, f_2, \dots \in \mathcal{H}_B$$

$$\psi = \sum_{\alpha} \lambda_{\alpha} e_{\alpha} \otimes f_{\alpha}$$

$$\lambda_{\alpha} \geq 0 , \quad \sum_{\alpha} \lambda_{\alpha}^2 = 1$$

Theorem:

$\psi \in \mathcal{H}_{\text{total}}$ is separable iff ONLY ONE $\lambda_{\alpha} \neq 0$

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Mixed states

Density matrix in \mathcal{H} : $\rho \geq 0$, $\text{Tr} \rho = 1$

A set of mixed states is CONVEX.

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Definition: (Werner 1989)

A bipartite state ρ is separable (classically correlated) iff

$$\rho = \sum_k p_k \rho_k^{(A)} \otimes \rho_k^{(B)}$$

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$\rho_k^{(A)}$ state in \mathcal{H}_A

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O_A & O_B observables in \mathcal{H}_A & \mathcal{H}_B

$$\rho = \sum_k p_k \rho_k^{(A)} \otimes \rho_k^{(B)}$$

$$\langle O_A \otimes O_B \rangle_\rho = \sum_k p_k \textcolor{red}{a_k} \textcolor{blue}{b_k}$$

$$\textcolor{red}{a_k} = \text{Tr} \left(O_A \cdot \rho_k^{(A)} \right) \quad \& \quad \textcolor{blue}{b_k} = \text{Tr} \left(O_B \cdot \rho_k^{(B)} \right)$$

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Separability problem

ρ — mixed state in $\mathcal{H}_{\text{total}} = \mathcal{H}_A \otimes \mathcal{H}_B$

Is ρ separable or entangled ?

To be separable is not a spectral property!!!

$$\rho = \sum_k p_k |\psi_k\rangle\langle\psi_k| \quad \text{spectral decomposition}$$

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Positive Partial Transpose (PPT)

$$\tau(X) = X^T$$

Definition:

A state ρ in $\mathcal{H}_A \otimes \mathcal{H}_B$ is PPT iff

$$(\mathbb{1} \otimes \tau)\rho \geq 0 .$$

Theorem: (Peres 1996)

Any separable state is PPT.

$$\rho = \sum_k p_k \rho_k^{(A)} \otimes \rho_k^{(B)}$$

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Separability vs. PPT

$$\text{SEPARABLE STATES} \subseteq \text{PPT STATES}$$

Theorem: (HHH 1996)

For

$$\mathbb{C}^2 \otimes \mathbb{C}^2, \mathbb{C}^2 \otimes \mathbb{C}^3, \mathbb{C}^3 \otimes \mathbb{C}^2$$

$$\text{SEPARABLE STATES} = \text{PPT STATES}$$

$$\text{If } \dim \mathcal{H}_A \cdot \dim \mathcal{H}_B > 6$$

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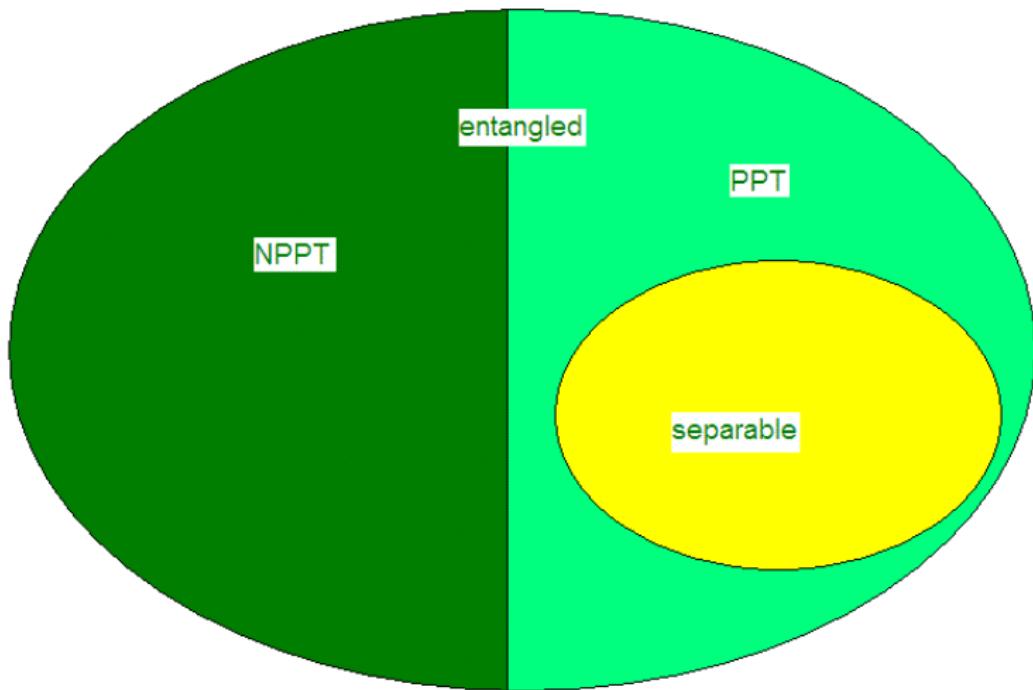
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3 \otimes 3 PPT but entangled

$$\rho = \frac{1}{1+8a} \begin{bmatrix} a & \cdot & \cdot & \cdot & a & \cdot & \cdot & \cdot & a \\ \cdot & a & \cdot \\ \cdot & \cdot & a & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & a & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & a & \cdot & \cdot & a & \cdot & \cdot & \cdot & a \\ \cdot & \cdot & \cdot & \cdot & \cdot & a & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \frac{1}{2}(1+a) & \cdot & \frac{1}{2}\sqrt{1-a^2} \\ \cdot & a & \cdot \\ a & \cdot & \cdot & \cdot & a & \cdot & \frac{1}{2}\sqrt{1-a^2} & \cdot & \frac{1}{2}(1+a) \end{bmatrix} \begin{array}{c} |00\rangle \\ |01\rangle \\ |02\rangle \\ |10\rangle \\ |11\rangle \\ |12\rangle \\ |20\rangle \\ |21\rangle \\ |22\rangle \end{array}$$

$\cdot \equiv 0, \quad a \in (0, 1)$



Separability vs. positive maps

\mathcal{A} & \mathcal{B} — C^* -algebras

$$\mathcal{A} \ni a \geq 0 \iff a = xx^*.$$

Definition:

A linear map

$$\Lambda : \mathcal{A} \longrightarrow \mathcal{B}$$

is positive iff

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Unfortunately, the structure of positive maps is poorly known!

17th Hilbert problem

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17th Hilbert problem



“Symmetry is one idea by which man through the ages has tried to comprehend and create order, beauty, and perfection.”

Symmetric states

$$\mathcal{H}_{\text{total}} = \mathcal{H}_A \otimes \mathcal{H}_B$$

G — compact Lie group

\mathfrak{D}_A & \mathfrak{D}_B unitary IRREPs

$$\mathfrak{D}_A(g) : \mathcal{H}_A \longrightarrow \mathcal{H}_A$$

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Example — Werner state

$$\mathcal{H}_A = \mathcal{H}_B = \mathbb{C}^d$$

$$G = U(d)$$

$\mathfrak{D}_A = \mathfrak{D}_B = \text{defining representation}$

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Werner state

$$(U(d) \otimes U(d))' = \{ \mathbb{I} \otimes \mathbb{I}, \mathbb{F} \}$$

$$\mathbb{F}(\psi \otimes \varphi) = \varphi \otimes \psi$$

$$\mathbb{F} = \sum_{i,j=1}^d e_{ij} \otimes e_{ji}$$

$$e_{ij} = |i\rangle\langle j|$$

Werner state

$$Q^0 = \frac{1}{2}(\mathbb{I}^{\otimes 2} + \mathbb{F}) , \quad Q^1 = \frac{1}{2}(\mathbb{I}^{\otimes 2} - \mathbb{F})$$

$$Q^0 + Q^1 = \mathbb{I}^{\otimes 2}$$

General form

$$\rho = q_0 \tilde{Q}^0 + q_1 \tilde{Q}^1 ; \quad \tilde{Q}^\alpha = Q^\alpha / \text{Tr } Q^\alpha ,$$

$$q_\alpha \geq 0 , \quad q_0 + q_1 = 1$$

Theorem

$$\rho - \text{separable} \iff \rho - \text{PPT} \iff q_0 \geq \frac{1}{2}$$

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SU(2) Symmetry

two particles \longleftrightarrow spin- $j_A \otimes$ spin- j_B $(j_B \geq j_A)$

$$\mathcal{H}_A = \mathbb{C}^{d_A} \quad \& \quad \mathcal{H}_B = \mathbb{C}^{d_B}$$

$$d_A = 2j_A + 1 \quad \& \quad d_B = 2j_B + 1$$

$$G = SU(2) \longrightarrow \mathfrak{D}^{(j_A)} \otimes \mathfrak{D}^{(j_B)}$$

$$\mathfrak{D}^{(j_A)} \otimes \mathfrak{D}^{(j_B)} = \bigoplus_{J=j_B-j_A}^{j_A+j_B} \mathfrak{D}^{(J)},$$

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$$\mathbb{C}^{2j_A+1} \longrightarrow |j_A, m_A\rangle ; \quad m_A = -j_A, \dots, +j_A$$

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$$\mathcal{H}_{\text{total}} = \mathbb{C}^{2j_A+1} \otimes \mathbb{C}^{2j_B+1} \longrightarrow |j_A, m_A\rangle \otimes |j_B, m_B\rangle$$

$$|j_A, m_A\rangle \otimes |j_B, m_B\rangle \longrightarrow |J, M\rangle$$

$$J = j_B - j_A, \dots, j_A + j_B$$

$$M = -J, \dots, +J$$

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$2j_B$ dimensional simplex

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$$\left[\mathfrak{D}^{(j_A)} \otimes \mathfrak{D}^{(j_B)}, Q^J \right] = 0$$

General form

$$\rho = \sum_{J=j_B-j_A}^{j_A+j_B} q_J \tilde{Q}^J,$$

$$q_J \geq 0, \quad \sum_J q_J = 1$$

$2j_B$ dimensional simplex

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SU(2) symmetry vs. separability

- $j_A = 1/2$ and arbitrary $j_B \geq j_A$

$$d_A \otimes d_B = 2 \otimes k ; \quad k = 2, 3, \dots$$

ρ is separable \iff ρ is PPT

- $j_A = 1$ and arbitrary integer $j_B \geq j_A$

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One may investigate other groups $G \subset U(d)$

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- symplectic group
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One may generalize construction of symmetric states to multipartite case

$$\mathcal{H}_{\text{total}} = \mathcal{H}_1 \otimes \dots \otimes \mathcal{H}_N$$

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Let us look for PPT states

It is very hard to characterize a set of separable states

SEPARABLE STATES \subseteq PPT STATES

Let us try to characterize a set of PPT states

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How to construct PPT states?

A state ρ in $\mathcal{H}_A \otimes \mathcal{H}_B$ is PPT iff

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PPT is easy to check!

But, it is not easy to construct ρ which is PPT!!!

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$$\mathcal{H}_A = \mathcal{H}_B = \mathbb{C}^d$$

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$$\rho = \sum_{i,j=0}^{d-1} a_{ij} e_{ij} \otimes e_{ij} + \sum_{i \neq j=0}^{d-1} b_{ij} e_{ii} \otimes e_{jj}$$

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Theorem: D.C. & A. Kossakowski, PRA (2006)

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New class in $d \otimes d$

Now we construct a wide class of states in $\mathcal{H}_{\text{total}} = \mathbb{C}^d \otimes \mathbb{C}^d$ which is based on symmetric decompositions of $\mathcal{H}_{\text{total}}$.

It generalizes the previous class invariant under $\mathbb{T}(d) \subset U(d)$.

Let us consider $d = 3$.

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Circulant decomposition of $\mathbb{C}^3 \otimes \mathbb{C}^3$

$$\mathbb{C}^3 \longleftrightarrow e_0, e_1, e_2$$

$$S : \mathbb{C}^3 \longrightarrow \mathbb{C}^3$$

$$S e_i = e_{i+1} \quad \text{mod } 3$$

$$\Sigma_0 = \text{span} \{e_0 \otimes e_0, e_1 \otimes e_1, e_2 \otimes e_2\} \in \mathbb{C}^3 \otimes \mathbb{C}^3$$

$$\Sigma_1 = (\mathbb{1} \otimes S) \Sigma_0$$

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Circulant state — construction

a , b , c 3×3 matrices

$$\begin{pmatrix} a_{00} & a_{01} & a_{02} \\ a_{10} & a_{11} & a_{12} \\ a_{20} & a_{21} & a_{22} \end{pmatrix}, \quad \begin{pmatrix} b_{00} & b_{01} & b_{02} \\ b_{10} & b_{11} & b_{12} \\ b_{20} & b_{21} & b_{22} \end{pmatrix}, \quad \begin{pmatrix} c_{00} & c_{01} & c_{02} \\ c_{10} & c_{11} & c_{12} \\ c_{20} & c_{21} & c_{22} \end{pmatrix}.$$

$$a \geq 0, \quad b \geq 0, \quad c \geq 0$$

$$\Sigma_0 = \text{span} \{e_0 \otimes e_0, e_1 \otimes e_1, e_2 \otimes e_2\}$$

$$\rho_0 = \begin{pmatrix} a_{00} & \cdot & \cdot & | & a_{01} & \cdot & | & \cdot & \cdot & a_{02} \\ \cdot & \cdot & \cdot & | & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & | & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ \hline \cdot & \cdot & \cdot & | & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ a_{10} & \cdot & \cdot & | & a_{11} & \cdot & | & \cdot & \cdot & a_{12} \\ \cdot & \cdot & \cdot & | & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ \hline \cdot & \cdot & \cdot & | & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & | & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ a_{20} & \cdot & \cdot & | & a_{21} & \cdot & | & \cdot & \cdot & a_{22} \end{pmatrix}.$$

$$\Sigma_1 = \text{span} \{e_0 \otimes e_1, e_1 \otimes e_2, e_2 \otimes e_0\}$$

$$\rho_1 = \left(\begin{array}{ccc|ccc|ccc} \cdot & \cdot \\ \cdot & b_{00} & \cdot & \cdot & \cdot & b_{01} & b_{02} & \cdot & \cdot \\ \cdot & \cdot \\ \hline \cdot & \cdot \\ \cdot & \cdot \\ \cdot & b_{10} & \cdot & \cdot & \cdot & b_{11} & b_{12} & \cdot & \cdot \\ \hline \cdot & b_{20} & \cdot & \cdot & \cdot & b_{21} & b_{22} & \cdot & \cdot \\ \cdot & \cdot \end{array} \right).$$

$$\Sigma_2 = \text{span} \{e_0 \otimes e_2, e_1 \otimes e_0, e_2 \otimes e_1\}$$

$$\rho_2 = \left(\begin{array}{ccc|ccc|ccc} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot & c_{00} & c_{01} & \cdot & \cdot & \cdot & c_{02} & \cdot \\ \hline \cdot & \cdot & c_{10} & c_{11} & \cdot & \cdot & \cdot & c_{12} & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \\ \hline \cdot & \cdot & c_{20} & c_{21} & \cdot & \cdot & \cdot & c_{22} & \cdot \\ \cdot & \cdot \end{array} \right).$$

Circulant state

$$\rho = \rho_0 + \rho_1 + \rho_2$$

$$\rho = \begin{pmatrix} a_{00} & \cdot & \cdot & | & \cdot & a_{01} & \cdot & | & \cdot & \cdot & a_{02} \\ \cdot & b_{00} & \cdot & | & \cdot & \cdot & b_{01} & | & b_{02} & \cdot & \cdot \\ \cdot & \cdot & c_{00} & | & c_{01} & \cdot & \cdot & | & \cdot & c_{02} & \cdot \\ \cdot & \cdot & c_{10} & | & c_{11} & \cdot & \cdot & | & \cdot & c_{12} & \cdot \\ a_{10} & \cdot & \cdot & | & \cdot & a_{11} & \cdot & | & \cdot & \cdot & a_{12} \\ \cdot & b_{10} & \cdot & | & \cdot & \cdot & b_{11} & | & b_{12} & \cdot & \cdot \\ \cdot & b_{20} & \cdot & | & \cdot & \cdot & b_{21} & | & b_{22} & \cdot & \cdot \\ \cdot & \cdot & c_{20} & | & c_{21} & \cdot & \cdot & | & \cdot & c_{22} & \cdot \\ a_{20} & \cdot & \cdot & | & \cdot & a_{21} & \cdot & | & \cdot & \cdot & a_{22} \end{pmatrix}.$$

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$$\rho_{\text{old}} = \begin{bmatrix} a_{00} & \cdot & \cdot & | & \cdot & a_{01} & \cdot & | & \cdot & \cdot & a_{02} \\ \cdot & b_{01} & \cdot & | & \cdot & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ \cdot & \cdot & b_{02} & | & \cdot & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ \hline \cdot & \cdot & \cdot & | & b_{10} & \cdot & \cdot & | & \cdot & \cdot & \cdot \\ a_{10} & \cdot & \cdot & | & \cdot & a_{11} & \cdot & | & \cdot & \cdot & a_{12} \\ \cdot & \cdot & \cdot & | & \cdot & \cdot & b_{12} & | & \cdot & \cdot & \cdot \\ \hline \cdot & \cdot & \cdot & | & \cdot & \cdot & \cdot & | & b_{20} & \cdot & \cdot \\ \cdot & \cdot & \cdot & | & \cdot & \cdot & \cdot & | & \cdot & b_{21} & \cdot \\ a_{20} & \cdot & \cdot & | & \cdot & a_{21} & \cdot & | & \cdot & \cdot & a_{22} \end{bmatrix}$$

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Partial transposition $\rho^\tau = (\mathbb{1} \otimes \tau)\rho$ — a MIRACLE!

$$\rho^\tau = \begin{pmatrix} \tilde{a}_{00} & \cdot & \cdot & | & \cdot & \tilde{a}_{01} & \cdot & | & \cdot & \tilde{a}_{02} & \cdot \\ \cdot & \tilde{b}_{00} & \cdot & | & \tilde{b}_{01} & \cdot & \cdot & | & \cdot & \cdot & \tilde{b}_{02} \\ \cdot & \cdot & \tilde{c}_{00} & | & \cdot & \tilde{c}_{01} & \cdot & | & \tilde{c}_{02} & \cdot & \cdot \\ \hline \cdot & \tilde{b}_{10} & \cdot & | & \tilde{b}_{11} & \cdot & \cdot & | & \cdot & \cdot & \tilde{b}_{12} \\ \cdot & \cdot & \tilde{c}_{10} & | & \cdot & \tilde{c}_{11} & \cdot & | & \tilde{c}_{12} & \cdot & \cdot \\ \tilde{a}_{10} & \cdot & \cdot & | & \cdot & \cdot & \tilde{a}_{11} & | & \cdot & \tilde{a}_{12} & \cdot \\ \hline \cdot & \cdot & \tilde{c}_{20} & | & \cdot & \tilde{c}_{21} & \cdot & | & \tilde{c}_{22} & \cdot & \cdot \\ \tilde{a}_{20} & \cdot & \cdot & | & \cdot & \cdot & \tilde{a}_{21} & | & \cdot & \tilde{a}_{22} & \cdot \\ \cdot & \tilde{b}_{20} & \cdot & | & \tilde{b}_{21} & \cdot & \cdot & | & \cdot & \cdot & \tilde{b}_{22} \end{pmatrix}$$

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Circulant decomposition of $\mathbb{C}^3 \otimes \mathbb{C}^3$

$$\tilde{\Sigma}_0 = \text{span} \{e_0 \otimes e_0, e_1 \otimes e_2, e_2 \otimes e_1\}$$

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$$\tilde{\Sigma}_1 = \text{span} \{e_0 \otimes e_1, e_1 \otimes e_0, e_2 \otimes e_2\}$$

$$\tilde{\rho}_1 = \left(\begin{array}{ccc|ccc|ccc} \cdot & \cdot \\ \cdot & \tilde{b}_{00} & \cdot & \tilde{b}_{01} & \cdot & \cdot & \cdot & \cdot & \tilde{b}_{02} \\ \cdot & \cdot \\ \hline \cdot & \tilde{b}_{10} & \cdot & \tilde{b}_{11} & \cdot & \cdot & \cdot & \cdot & \tilde{b}_{12} \\ \cdot & \cdot \\ \cdot & \cdot \\ \hline \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \tilde{b}_{20} & \cdot & \tilde{b}_{21} & \cdot & \cdot & \cdot & \cdot & \tilde{b}_{22} \end{array} \right)$$

$$\tilde{\Sigma}_2 = \text{span} \{e_0 \otimes e_2, e_1 \otimes e_1, e_2 \otimes e_0\}$$

$$\tilde{\rho}_2 = \left(\begin{array}{ccc|ccc|ccc} \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot & \tilde{c}_{00} & \cdot & \tilde{c}_{01} & \cdot & \tilde{c}_{02} & \cdot & \cdot & \cdot \\ \hline \cdot & \cdot \\ \cdot & \cdot & \tilde{c}_{10} & \cdot & \tilde{c}_{11} & \cdot & \tilde{c}_{12} & \cdot & \cdot & \cdot \\ \cdot & \cdot \\ \hline \cdot & \cdot & \tilde{c}_{20} & \cdot & \tilde{c}_{21} & \cdot & \tilde{c}_{22} & \cdot & \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \end{array} \right)$$

$$a, b, c \longrightarrow \tilde{a}, \tilde{b}, \tilde{c}$$

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Theorem

Circulant state ρ is PPT iff

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π & $\tilde{\pi}$ are “complementary”

$$\pi(i) + \tilde{\pi}(i) = 3 \quad i = 1, 2$$

$$\Pi = \begin{pmatrix} 1 & \cdot & \cdot \\ \cdot & 1 & \cdot \\ \cdot & \cdot & 1 \end{pmatrix}, \quad \tilde{\Pi} = \begin{pmatrix} 1 & \cdot & \cdot \\ \cdot & \cdot & 1 \\ \cdot & 1 & \cdot \end{pmatrix}$$

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This construction works for any d !!!

There are $(d - 1)!$ circulant decompositions labeled by permutations $\pi \in S_{d-1}$

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- Werner state
- Isotropic state
- DiVincenzo *et al.* BE state
- Størmer state
- Ha state
- ... and many others

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Open problems

- separability
- “circulant” entanglement witnesses
- “circulant” positive maps
- generalization to $d \otimes d \otimes \dots \otimes d$
- do there exist other characteristic decompositions?